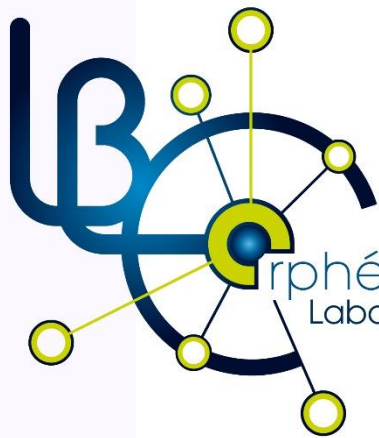


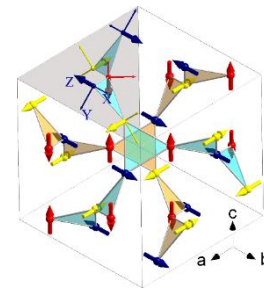
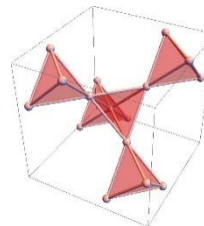
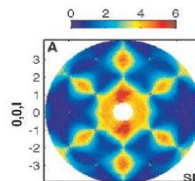
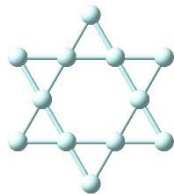


raphée
Laboratoire Léon Brillouin

SQL workshop, Budapest 2025

High entropy materials as a platform to explore frustrated magnets

F. Damay



ICMMO

F. Vayer, C. Decorse, D. Bérardan, N. Dragoe
(powder samples synthesis and characterizations)

LLB

S. Petit
(INS, modelling)

ILL

C. Colin (D1B)
S. Rols (Panther)

PSI

J. Embs (Focus)

Institut Néel

E. Lhotel
(LT magnetisation measurements)

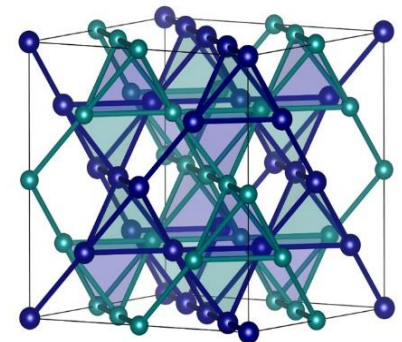




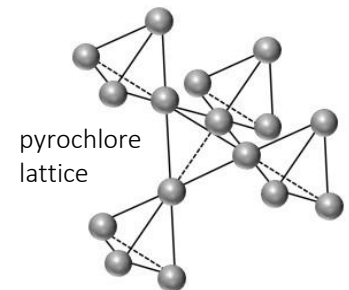
limited to A^{3+}/B^{4+}

site B (under pressure)

Quite a flexible structure!
Many compositions possible...
...but most have been investigated



A and B pyrochlore lattices



Disorder – an original tool to design new materials

In 2004, first high-entropy alloys (at least 5 cations on a given crystallographic site)

Free enthalpy of formation

$$\Delta G = \Delta H - T\Delta S$$

$$\Delta H < 0$$

Classical route

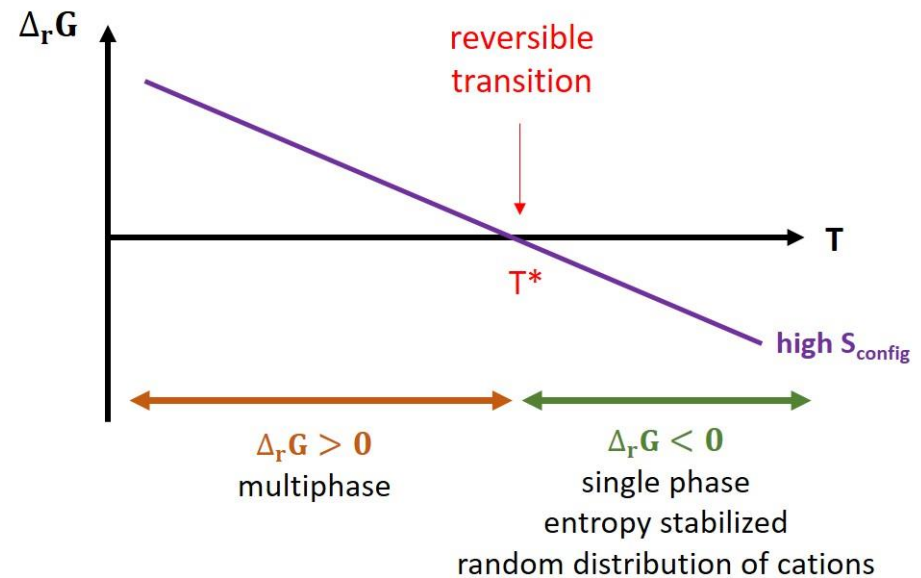
$$\Delta H > 0$$

Classical route fails
Stabilisation possible
through the $T\Delta S$ term
($\Delta G < 0$ at high T ,
METASTABLE at RT)

High Entropy compounds ($\Delta S_{\text{config}} = R \ln(N) \geq 1.61R$)

Entropy Stabilized compounds ($T\Delta S > \Delta H$)

- ★ disorder(s) and cocktail effect
- ★ stabilization of unusual geometries/valencies/structures



Very large range of
NEW compounds
NEW unexpected
properties

Design of an entropy-stabilized compounds with a pyrochlore structure

Some criteria

- preserve the A lattice, HEOx on the B site
- not the same structure/coordination in the parent compounds
- $r_{(avg)B}$ in the pyrochlore structure stability range
- equimolar mixture on B
- *(isovalent and non-magnetic cations to start with)*

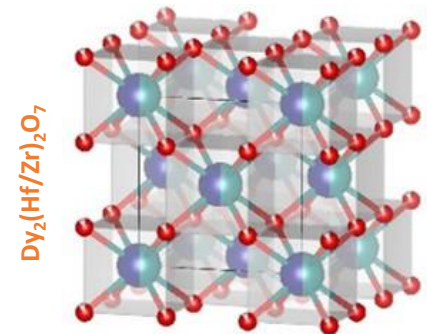
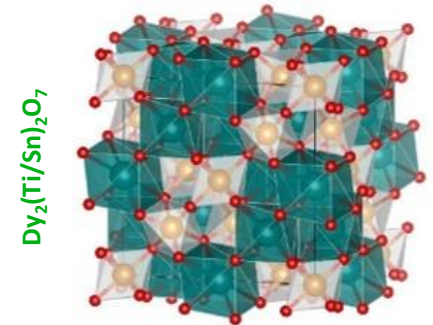
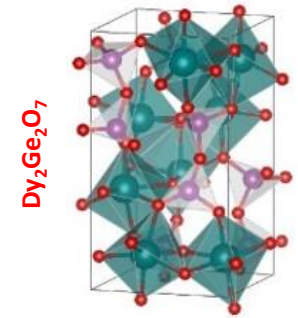
Parent compound	r_B	r_A/r_B	Structure	Magnetic behaviour
$Dy_2Ge_2O_7$	0,53	1,94	tetragonal or pyrochlore*	- spin ice
$Dy_2Ti_2O_7$	0,605	1,70	pyrochlore	spin ice
$Dy_2Sn_2O_7$	0,69	1,49	pyrochlore	spin ice
$Dy_2Hf_2O_7$	0,71	1,45	fluorite	-
$Dy_2Zr_2O_7$	0,72	1,43	fluorite	-

$Dy_2(TiZrHfGeSn)_2O_7$
(equimolar mixture on B site)

$r_{(avg)B}$
0,651 1,58

?

?



Synthesis at sufficiently high T (1600 °C) + quenching

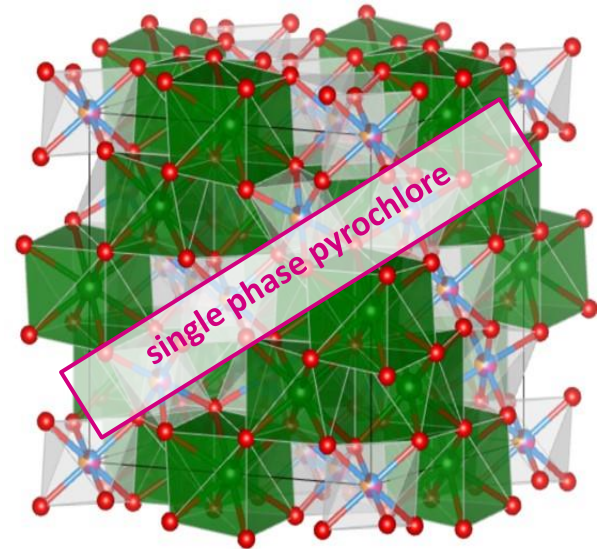


1st entropy stabilized pyrochlore : $\text{Dy}_2(\text{TiZrHfGeSn})_2\text{O}_7$

- single phase compound
- reversible transition between single phase HT and multiple phases at LT proves the entropy stabilization

$\text{A}_2\text{B}_2\text{O}_7$	a (Å)	T _{freeze} (K)	Ground state
$\text{Dy}_2\text{Ge}_2\text{O}_7$	9,93	0,83	spin ice
$\text{Dy}_2\text{Ti}_2\text{O}_7$	10,10	1,2	spin ice
$\text{Dy}_2(\text{TiZrHfGeSn})_2\text{O}_7$	10,30	?	? in progress!
$\text{Dy}_2\text{Sn}_2\text{O}_7$	10,40	1,2	spin ice

Robustness of the spin ice ground state in Dy pyrochlores



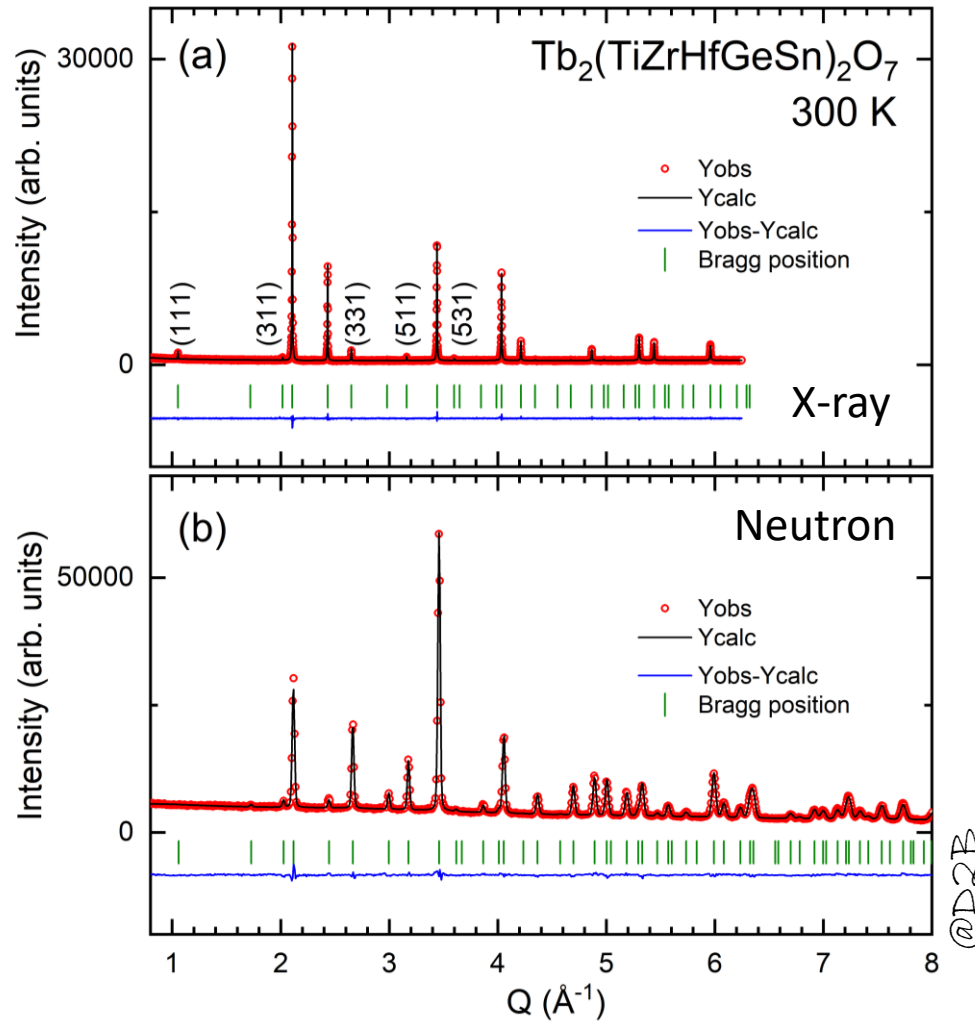
This new route offers rich perspectives for the design of original materials

- possibility to extend to all REs by changing the cations mixture
- multiplication of possibilities using mixtures of allivalent cations



New hints on the enigmatic ground state of $\text{Tb}_2\text{Ti}_2\text{O}_7$?

Crystal structure characterizations

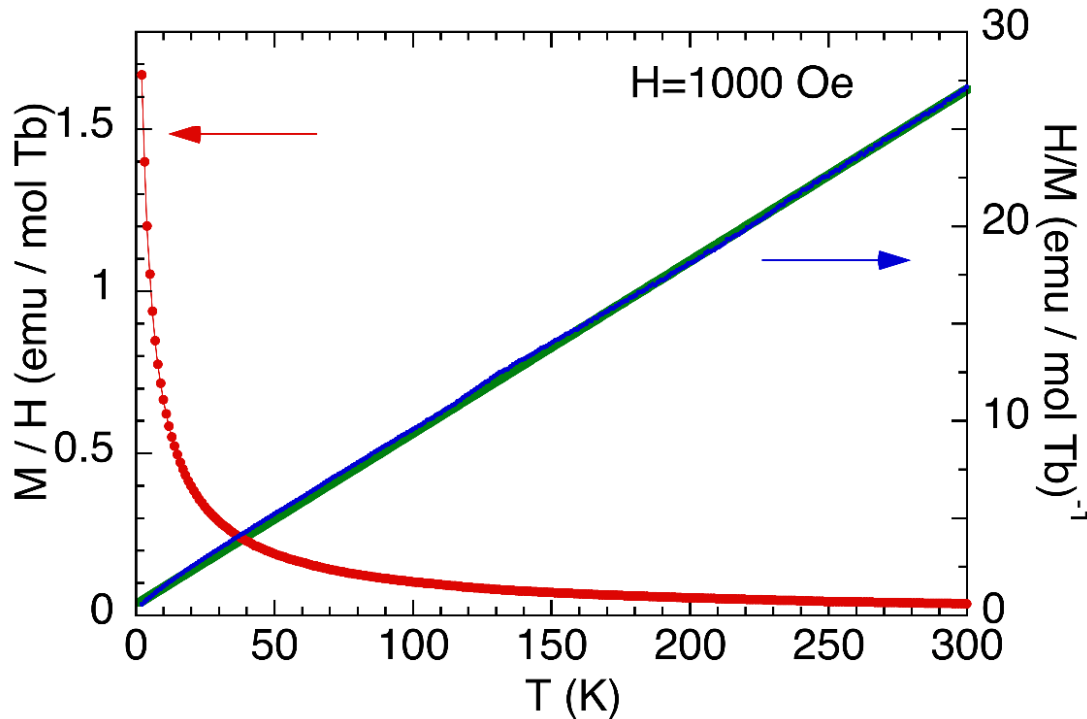


Tb on the A-site of the pyrochlore structure

Disorder introduced on the B site with 5 non-magnetic cations (Ti^{4+} , Zr^{4+} , Hf^{4+} , Ge^{4+} and Sn^{4+})

- **Single phase cubic $Fd-3m$**
 $a = 10.3287(10) \text{ \AA}$ at RT
 Confirmed stoichiometry
- **Debye-Waller** parameter large on O ($> 1 \text{ \AA}^2$) - position disorder
- **Antisite** (RE $\leftrightarrow$ B) disorder possible (as b for Tb, Hf, Ti are similar)
- **No broadening**, or short range structural disorder signature

Magnetization measurements

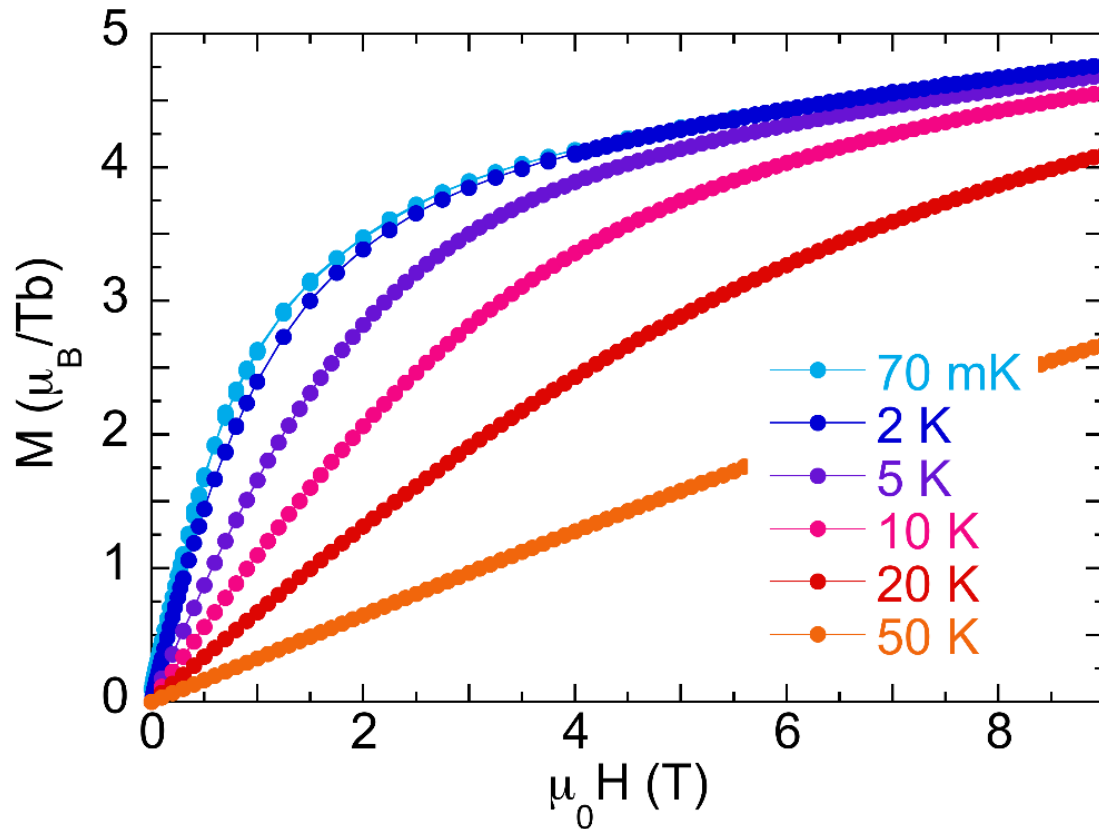


- No ordering down to 2 K
Almost linear behaviour above 20 K, like in other Tb pyrochlores
- $\theta_{\text{CW}} = -6.4 \text{ K}$ ($\approx \text{Tb}_2\text{ScNbO}_7$), $M_{\text{eff}} = 9.5 \mu_{\text{B}}$ ($gJ = 9 \mu_{\text{B}}$ for Tb^{3+})

$$\theta_{\text{CW}} = \theta_0 + \theta_{\text{CW}}^{\text{CEF}}, \text{ if } \theta_{\text{CW}}^{\text{CEF}} \approx -6 \text{ K}, \theta_0 \approx -0.5 \text{ K}$$

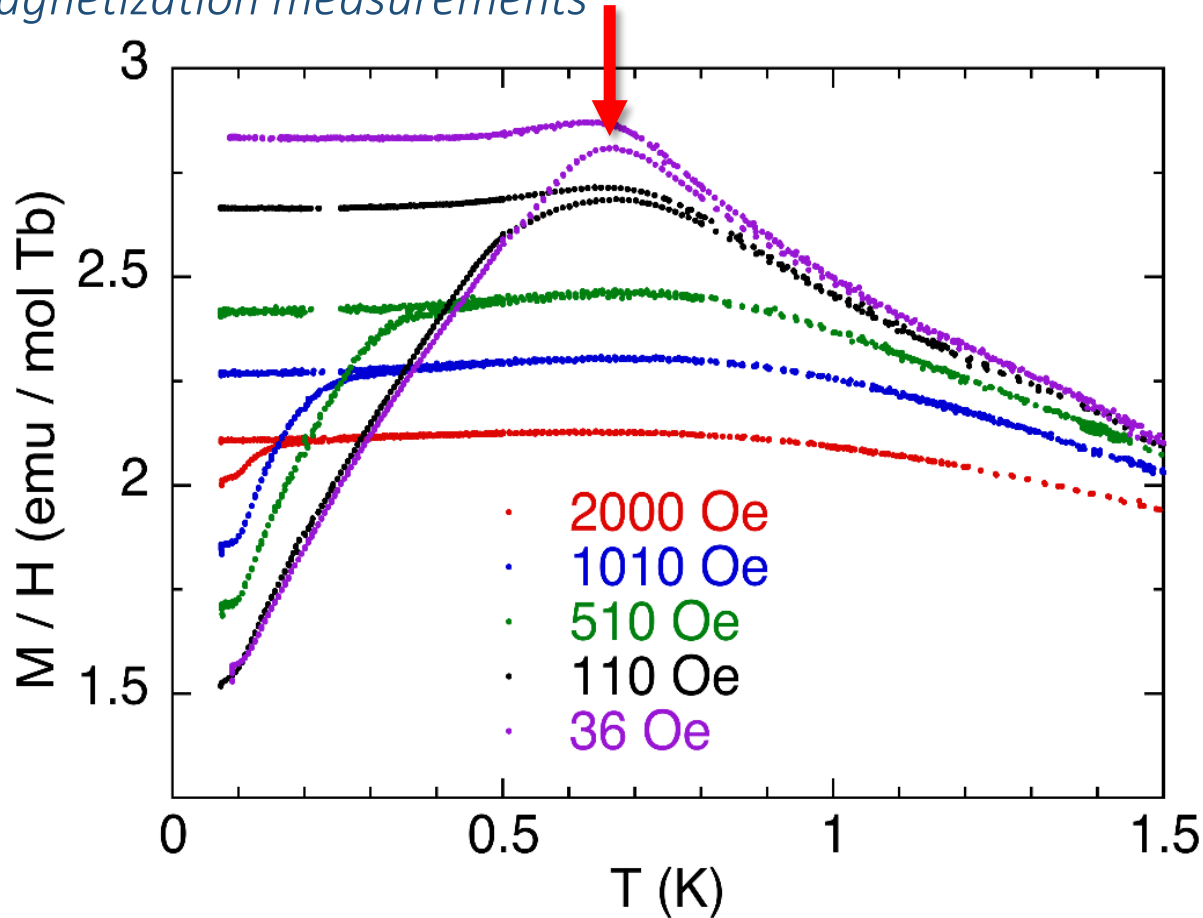
Weak AFM interactions

Magnetization measurements



- No saturation in 70 mK, 9 T
- $M_{\text{max}} \approx 4.75 \mu_B/\text{Tb}$, like $\text{Tb}_2\text{Ti}_2\text{O}_7$ or $\text{Tb}_2\text{Nb}_2\text{O}_7$

Zfc-fc magnetization measurements

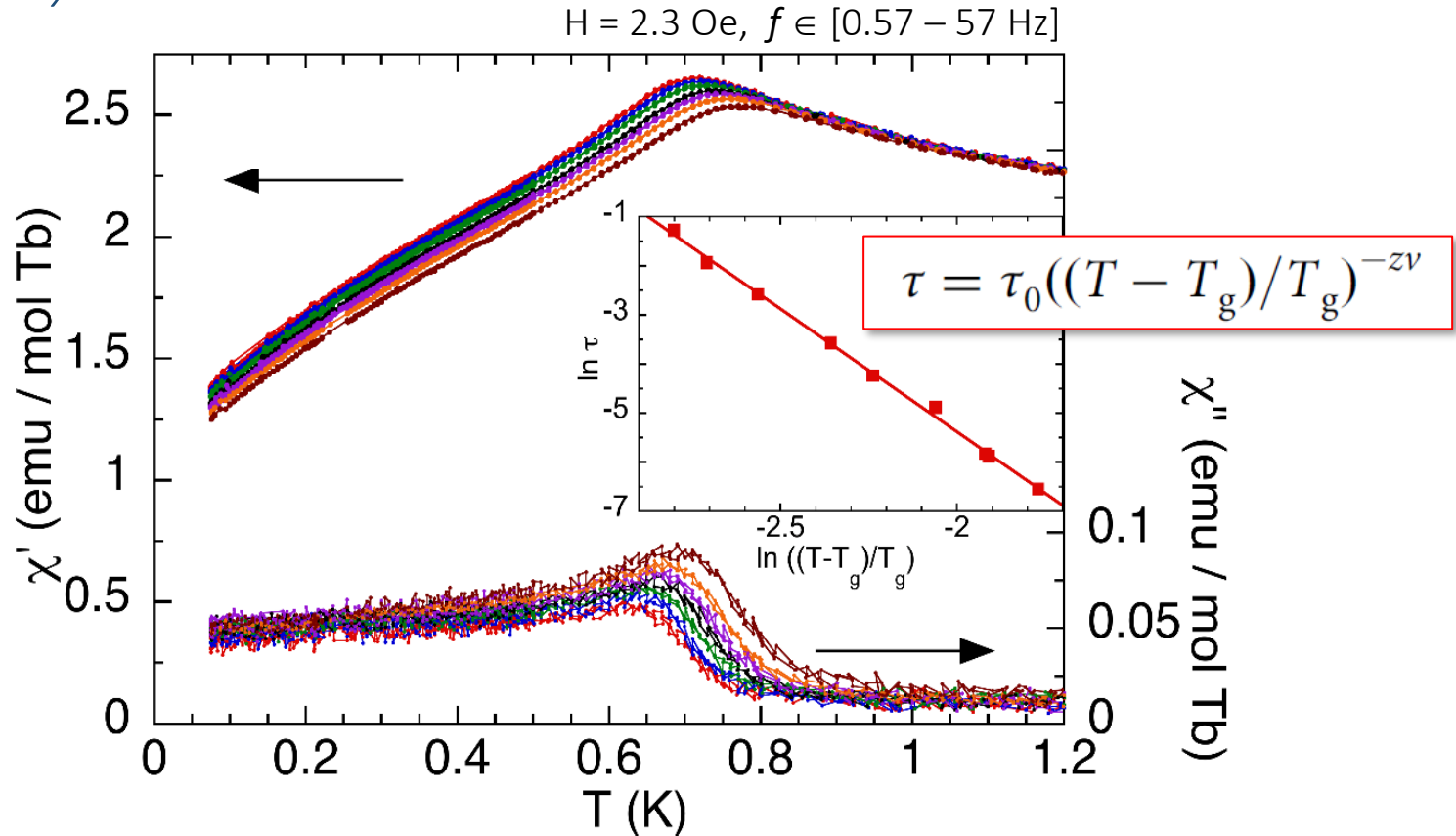


- Maximum at $T \approx 0.65$ K, irreversibility ≈ 0.8 K
- Broadening when H increases, irreversibility suppressed for $H > 0.5$ T

—→ Spin freezing, $r \approx 50\%$ ($r = 1 - M_{\text{ZFC}}/M_{\text{FC}}$)

like $\text{Tb}_2\text{ScNbO}_7$

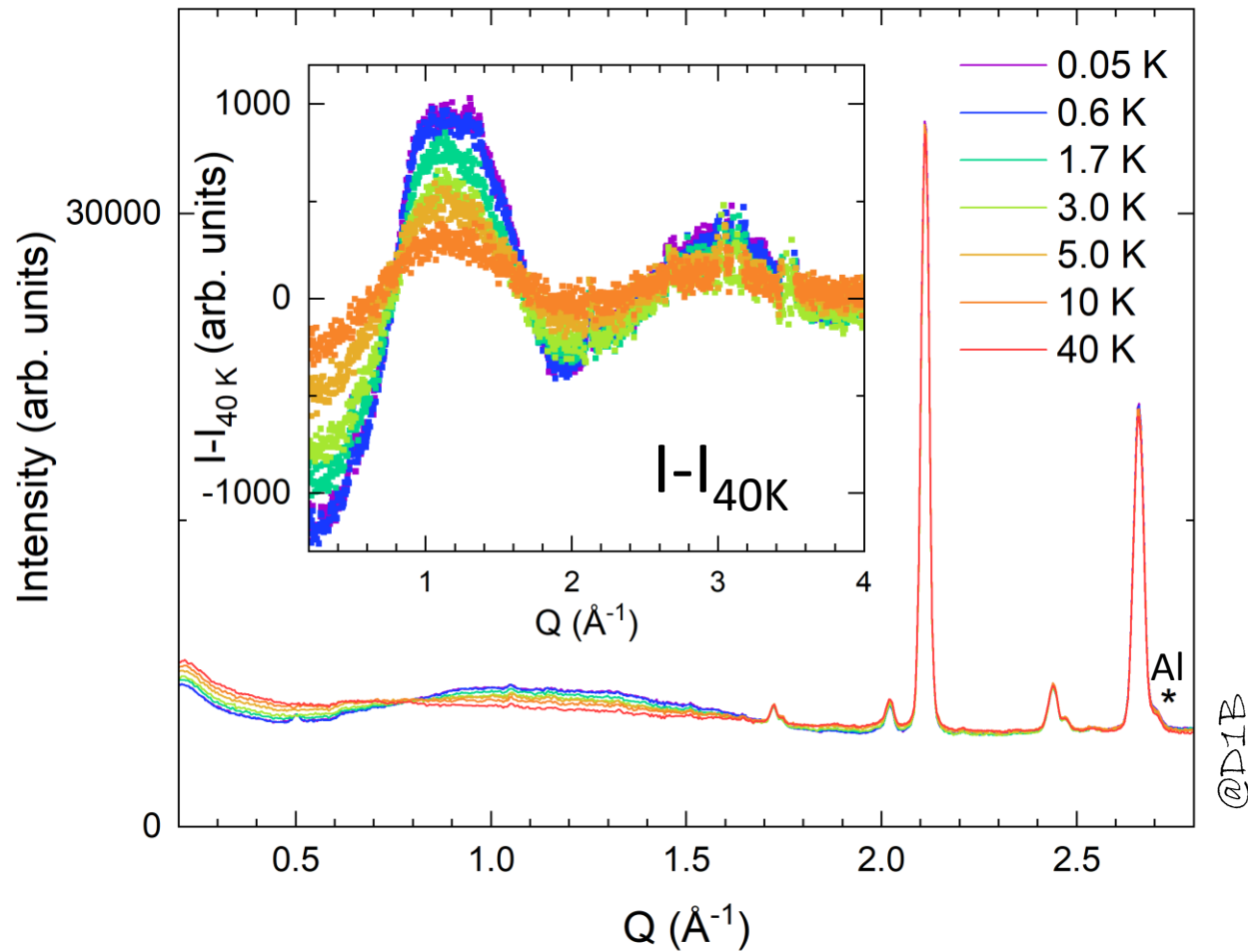
Ac susceptibility measurements



- Frequency dependence of χ' and χ''
- Described by a scaling law
- Similar parameters in other disordered Tb pyrochlores ($\text{Tb}_2\text{Hf}_2\text{O}_7$ and $\text{Tb}_2\text{ScNbO}_7$)

$$\tau_0 = 2.10^{-7} \pm 5.10^{-8} \text{ s}, T_g = 0.675 \pm 0.05 \text{ K and } zv = 5 \pm 1$$

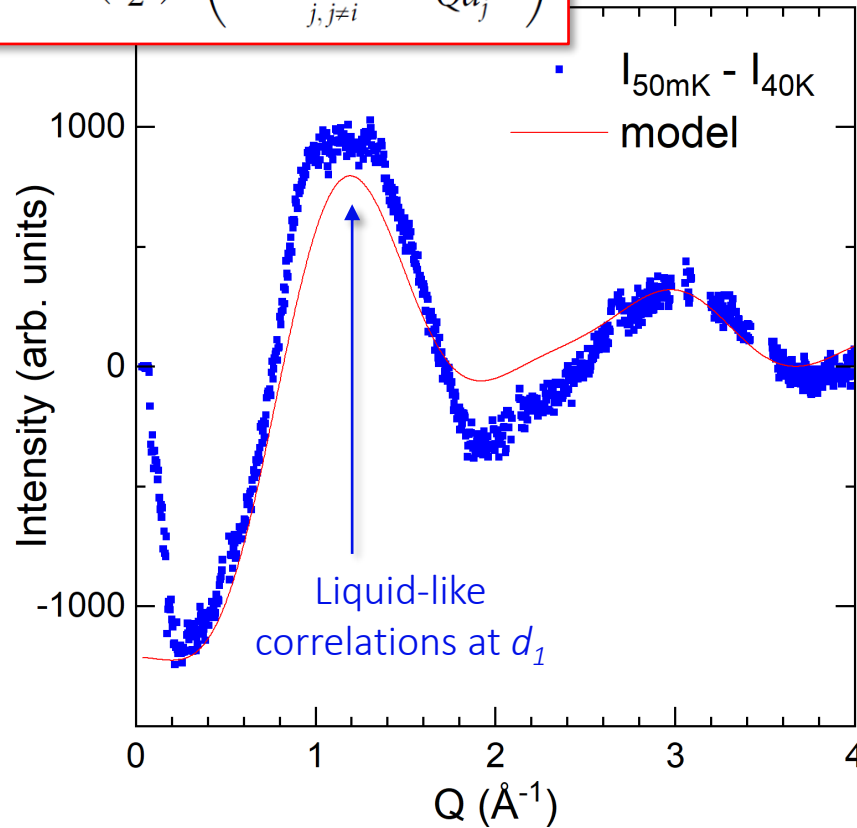
Neutron powder diffraction



- No long range ordering down to 50 mK
- Strong diffuse scattering signal

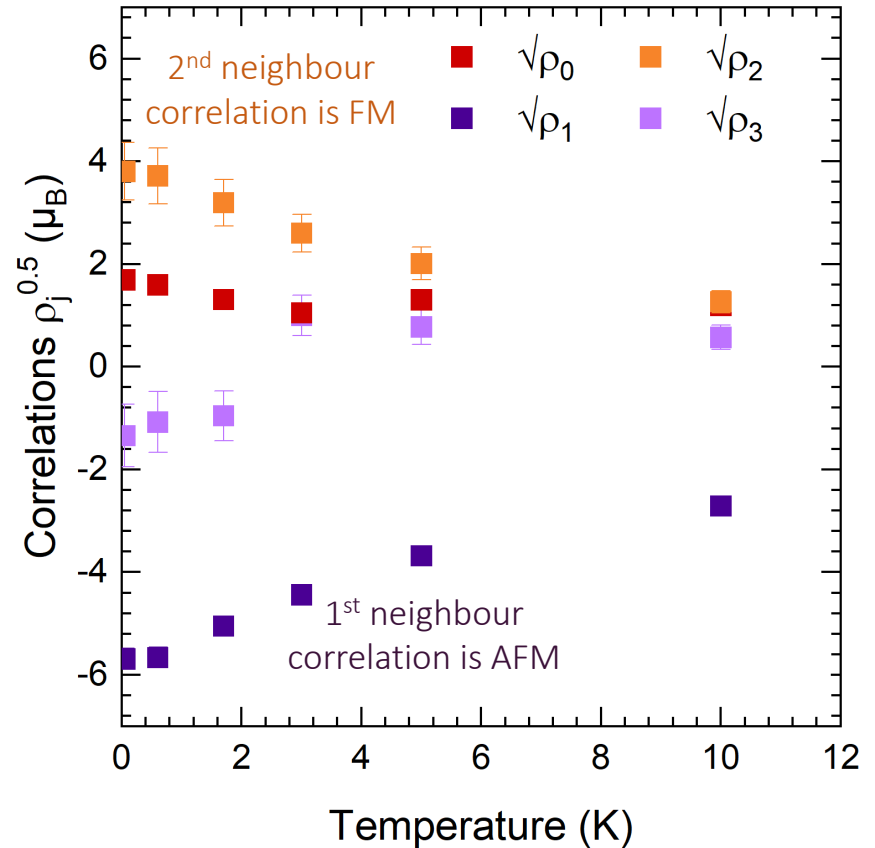
Correlated spin modelling of the magnetic diffuse scattering

$$I_M = N \left(\frac{\gamma r_0}{2} \right)^2 \left(\rho_0 + \sum_{j,j \neq i} \rho_j \frac{\sin(Qd_j)}{Qd_j} \right)$$



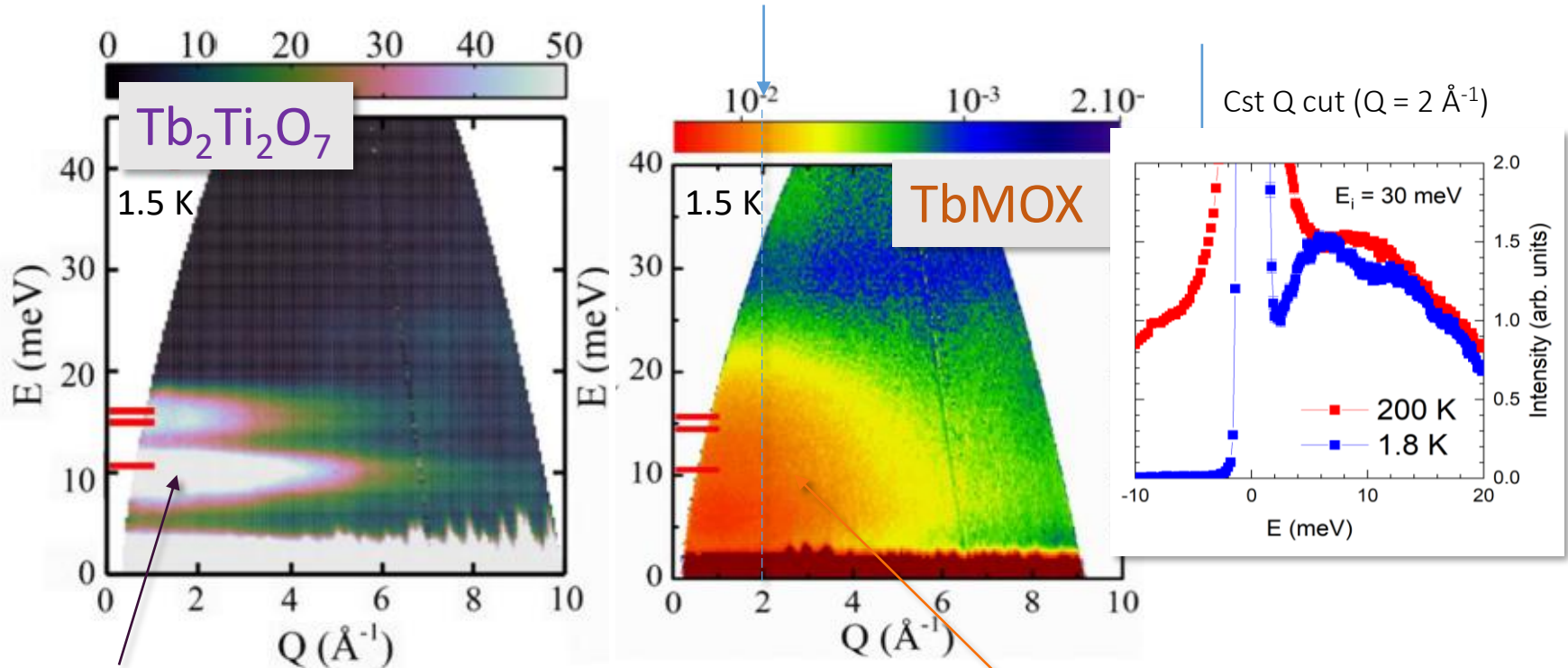
General shape of the diffuse scattering reproduced by the modelling

Evolution with T of the correlated moments



Similar diffuse scattering features to $\text{Tb}_2\text{Ti}_2\text{O}_7$

Probing CEF excitations by inelastic neutron scattering (0 – 40 meV)



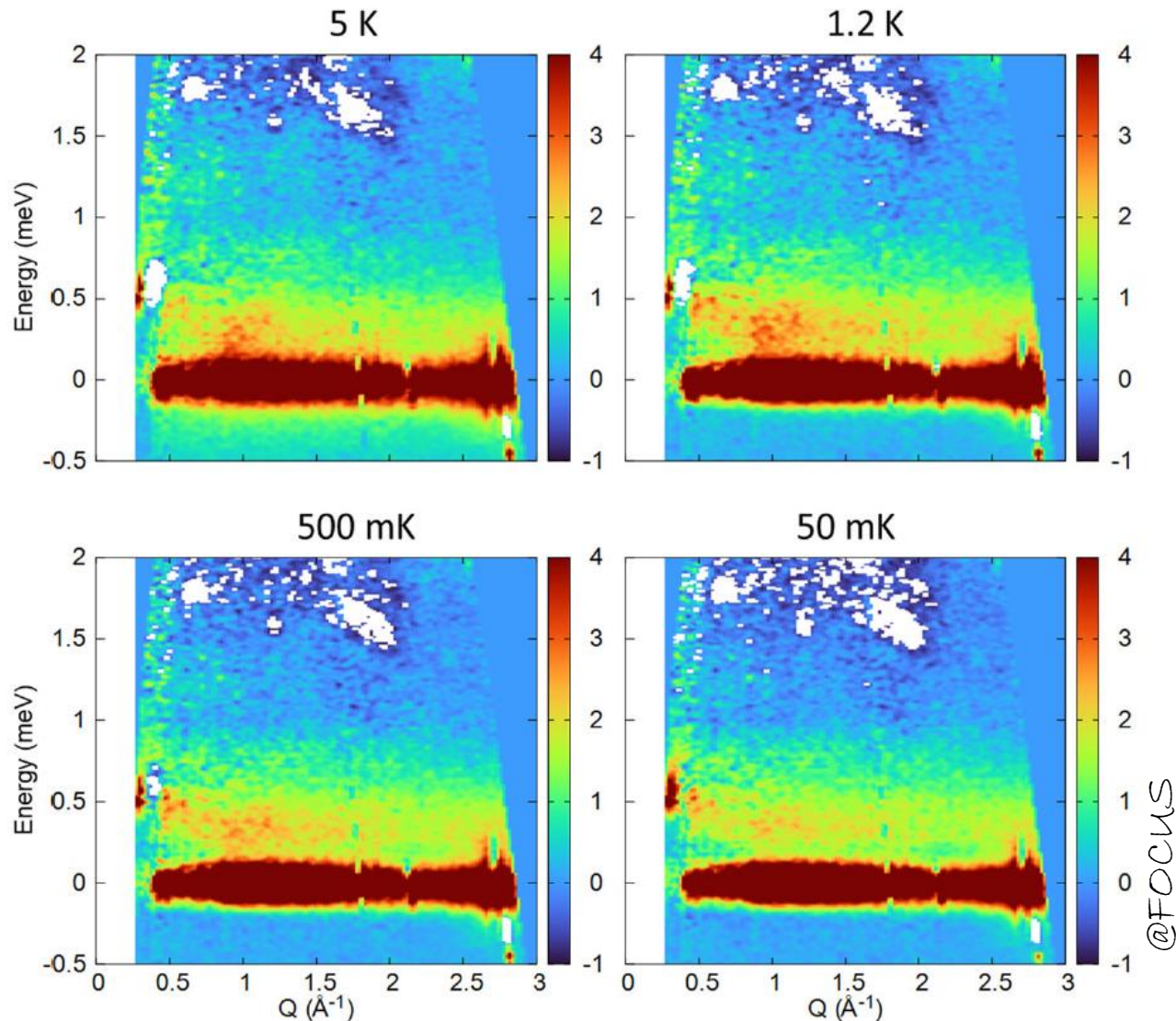
@Panther

Several discrete and flat levels below 20 meV*

*Ruminy, Magnetoelastic effects in rare earth Titanate Pyrochlores, Thèse de Doctorat, ETH Zurich, (2016).

Broad signal between 4 and 40 meV
Reflects the distribution of oxygen positions around Tb ions

Probing CEF excitations by inelastic neutron scattering (0 – 2 meV)



At low E, a quasi
 elastic signal
 transforms into an
 inelastic flat level at
 $E \approx 0.37 \text{ meV}$
 Strikingly narrow
 level ($\Gamma = 0.25 \text{ meV}$)

@FOCUS

Modelling (CEF part)

CEF Hamiltonian

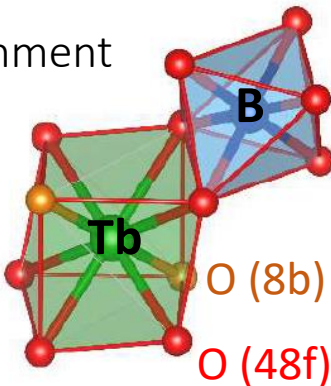
$$H_{\text{CEF}} = \sum_{n,m} B_{n,m} O_{n,m}$$

Stevens parameters $\nearrow$ $B_{n,m}$ $O_{n,m}$ $\nearrow$ Stevens operators

For a pyrochlore A site (D_{3d} symmetry) : B_{20} , B_{40} , B_{43} , B_{60} , B_{63} , B_{66} only are non zero

Simple approach to estimate B_{nm} : **the point charge model**

Tb environment
 D_{3d}



To model B-site entropy : oxygen coordinates are shifted by an isotropic shift parameter δ , away from their actual values.

$\delta \approx$ Debye-Waller parameter

Importantly : random δ on the 8 oxygen atoms introduce a **symmetry breaking** of the local D_{3d} symmetry.

All B_{nm} now allowed, D_{3d} symmetry fulfilled only in average

Using the point charge approximation, CEF energy levels and corresponding neutron scattering spectrum can be calculated vs δ

Modelling (magnetic interactions part)

Magnetic diffuse scattering implies magnetic interactions between Tb moments :

$$\mathcal{H}_{\text{mag}} \approx \sum_{\langle i,j \rangle} g^{zz} I_i^z I_j^z$$

Owing to the non-Kramers nature of the Tb^{3+} ions

$z \parallel \langle 111 \rangle$, g^{zz} the magnetic coupling

To keep the approach simple, a mean field approximation is used, involving a **local molecular field H**

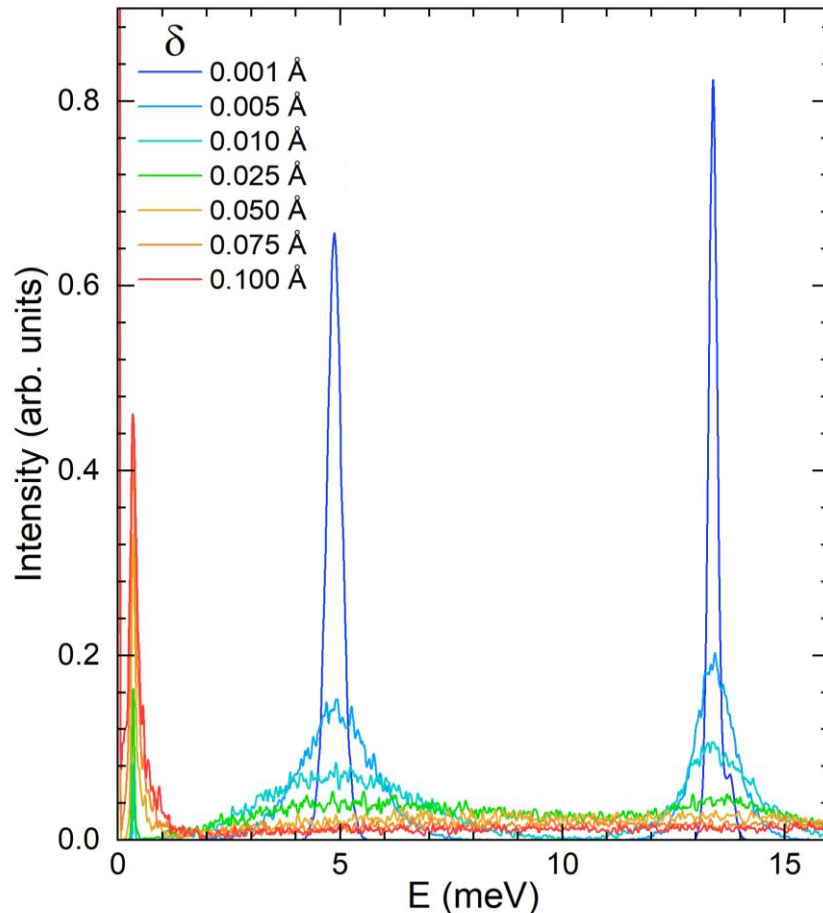
Modelling (CEF + magnetic interactions)

$$\mathcal{H} \approx \mathcal{H}_{\text{CEF}} + H I_i^z$$

This model only has two parameters, the **molecular field H** and the **oxygen position shift δ**

Modelling (CEF + magnetic interactions)

Calculations CEF spectra vs. δ for $H = 0.45$ K



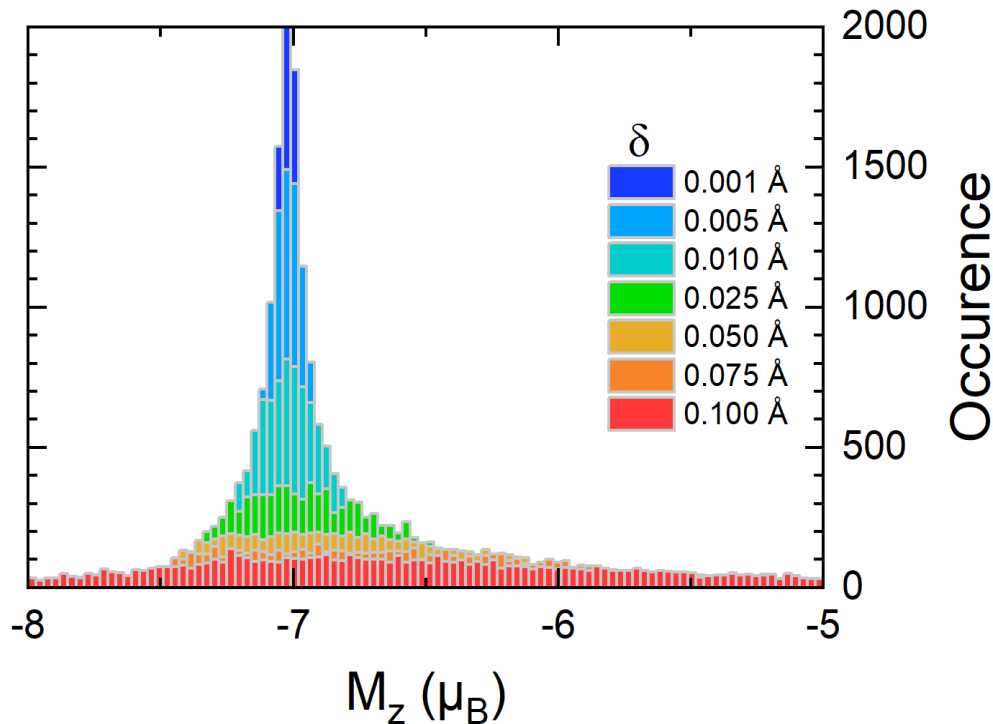
$H = 0.45$ K leads to a split ground state with a first level at 0.4 meV

- With increasing δ (disorder), clear broadening of the CEF levels at high energy
- Spectral weight of the low E mode (at 0.4 meV) is zero if $\delta = 0$
- Spectral weight of 0.4 meV mode increases with increasing δ , without broadening or E shift

Disorder must be present to visualize this mode!

Modelling (CEF + magnetic interactions)

Distribution of moment value vs. δ for $H = 0.45 \text{ K}$



The distribution of Tb^{3+} moment values broadens also with δ , but remains peaked around $7 \mu_B$ up to $\delta = 0.05 \text{ Å}$.

Nearly uniform distribution of local moments!

This simple modelling reproduces rather well experimental findings!

Can we go further....?

More complex modelling.....introducing quadrupole interactions

In $\text{Tb}_2\text{Ti}_2\text{O}_7$, quadrupolar interactions are a key ingredient to the understanding of the ground state

A simplified way to tackle quadrupolar interactions is to use a projection of the Hamiltonian on a subspace formed by the two states $\uparrow$ and $\downarrow$ of a pseudo spin $\frac{1}{2}$ (σ_i). This approach only works for the ground doublet, which is the one of interest here. At the mean field level :

$$\mathcal{H}' \approx \sum_i h \sigma_i^z + \sum_i h_{\perp} \sigma_i^x$$

Longitudinal molecular
field due to magnetic
interactions

Transverse effective field

$$h_{\perp} = (\bar{h}_{\perp} + \delta h_{\perp})$$

Quadrupole
-quadrupole
interactions

Net effect of random
deviations out of D_{3d}
symmetry

The z component of the pseudo spin is directly related to the local magnetic moment M_z

$$M_z = g_J \langle I_i^z \rangle \approx 7 \mu_B$$

More complex modelling.....introducing quadrupole interactions

Diagonalization of $\mathcal{H}'$ shows that there are two eigenstates separated by :

$$\Delta = \sqrt{h^2 + h_{\perp}^2}$$

Corresponding neutron spectrum :

$$S(\omega) \approx \mu^2 / 4 \left\{ \frac{h^2}{h^2 + h_{\perp}^2} \delta(\omega) + \frac{h_{\perp}^2}{h^2 + h_{\perp}^2} \delta(\omega - \Delta) \right\}$$

Complementary spectral weight

Elastic

Inelastic

Physical understanding of the low energy peak and its behaviour with δ

Peak at finite Δ whose spectral weight depends on $h_{\perp}$

More complex modelling.....introducing quadrupole interactions

Approximations

$$\Delta = \sqrt{h^2 + h_{\perp}^2} = 0.37 \text{ meV}$$

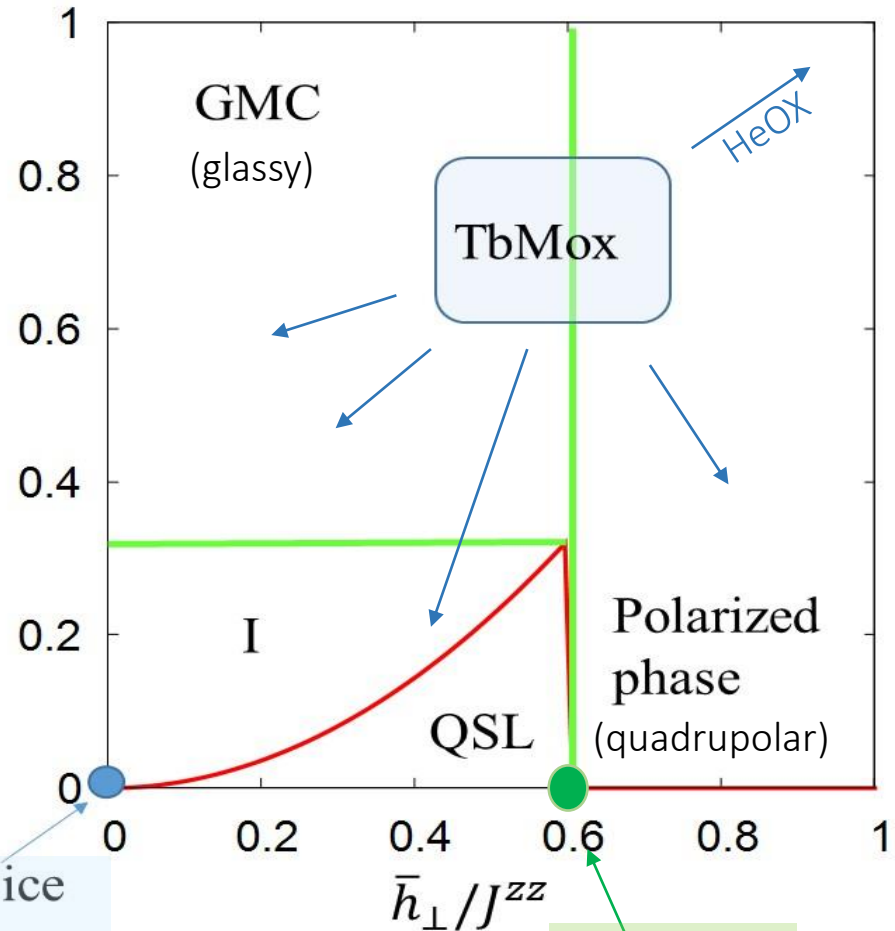
$$\delta h_{\perp} = 0.25 \text{ meV } (\approx \Gamma)$$

$$h \approx J^{zz} \approx 0.32 \text{ meV } (\sqrt{\rho_1} = 6 \mu_B)$$

$$\frac{\delta h_{\perp}}{J^{zz}} \sim 0.8, \frac{\bar{h}_{\perp}}{J^{zz}} \sim 0.6$$

Disorder and quadrupolar effects
of same order of magnitude

$$\frac{\delta h_{\perp}}{J^{zz}}$$



Spin ice

$$\bar{h}_{\perp} = \delta h_{\perp} = 0$$

(?) $\text{Tb}_2\text{Ti}_2\text{O}_7$

Conclusions

- A new, entropy-stabilized terbium pyrochlore!
- Like in $\text{Tb}_2\text{Ti}_2\text{O}_7$: AF interactions and a narrow low energy flat mode
- Compositional disorder leads to random shifts of oxygen atoms
- Disorder
 - does not impact Tb^{3+} magnetic moment
 - broadens high energy CEF levels
 - increases low energy excitation mode intensity
- High entropy provides a tunable and smooth disorder : promising to achieve QSL phases!

Thank you for your attention!

